

A retrospective look at coupled ice sheet–climate modeling

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Since its inception, numerical climate modeling has evolved along with available computer power. Limitations in computer power quickly led to distinct types of models, with relatively simple models capable of long integrations, versus complex models suitable for short-duration detailed snapshots. Recognizing these computing limitations, strategies to combine and enhance knowledge from the different model types were conceived, and Schneider and Dickinson (1974) were early proponents of interactive “hierarchies” of models. In that framework, numerical climate models of different complexities, ranging from energy balance models (EBMs) for long-term simulations, through zonal statistical-dynamic models and nowadays, Earth models of intermediate complexity (EMICs), to general circulation (or global climate) models (GCMs) for short-term weather details, are used in combination with each other. For instance, the GCM is used to determine key sensitivities (to orbital perturbations, for example), and then the EBM is tuned to have the same sensitivities. Knowledge and experience at each level of the hierarchy is applied interactively at other levels. Climate-model hierarchies have also been discussed by Henderson-Sellers and McGuffie (1987), Claussen et al. (2002) and Bartlein and Hostetler (2004), with Claussen et al. (2002) distinguishing between integration of components vs. detail of description, and proposing the term “spectrum” to avoid any suggestion that one hierarchical level is better than another (Fig. 1).

This paper briefly surveys how these ideas have found form over the last several decades, in the area of coupled ice sheet–climate modeling. To some extent the original concept of hierarchies has been realized, but mostly it has been adapted in different ways than originally envisioned, driven by the need to address the very

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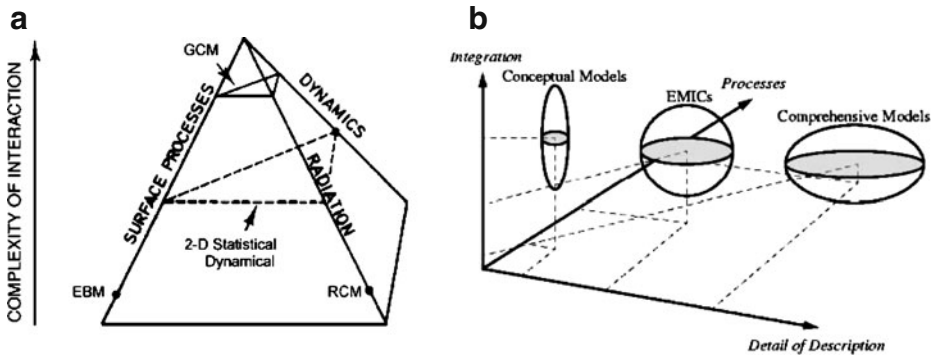


Fig. 1 **a** Climate-modeling pyramid, adapted from Henderson-Sellers and McGuffie (1987). **b** Spectrum of climate models, from Claussen et al. (2002)

different temporal and spatial scales of synoptic meteorology (days, $\sim 10^3$ km) and ice sheets (10^3 to 10^5 years, 10^3 km near their margins; Saltzman 1984, 1990). Techniques addressing the time-scale mismatch are described first, followed by a shorter discussion of spatial downscaling methods. The last section summarizes other types of hierarchical interactions and future directions.

1 Multiple time scales

Since the early studies of causes of the ice ages in the nineteenth and early twentieth centuries, a tight connection between ice cover and climate was realized, and the need to involve meteorologic forcing by orbital perturbations (Adh mar, Croll, Milankovitch, described in Imbrie and Imbrie 1979). As data and theories became more refined in the late 20th century, the need to represent climate with more robust models emerged. However, the mismatch between “fast” weather/climate variations and much slower ice-sheet variations is a fundamental problem for numerical modeling; on the one hand, there is a need to accurately simulate weather over an ice sheet to provide its annual surface mass balance, but doing so straightforwardly with full GCMs driving ice sheet models over million-year time scales would take prohibitively long even on today’s computers.

Nevertheless, long-term simulations are needed to address many ice-age questions. Continental-scale ice sheets typically vary on time scales of 10^4 years or more, and the fastest large-scale fluctuations such as the final collapse before interglacials still take several 1,000’s of years. These fluctuations occur repeatedly throughout ice ages of 10^6 to 10^7 years extent, such as Neoproterozoic Snowball Earth events (Hoffman and Shrag 2000), sudden growth of major Antarctic ice and subsequent oscillations in the Cenozoic, ~ 34 Ma to the present (Zachos et al. 2006, 2008), and Quaternary Northern Hemispheric glacial/interglacial cycles of the last ~ 3 Myear (Imbrie et al. 1993).

Progress in understanding these events requires the ability to drive prognostic models of the long-term components (ice sheets, sediment and bedrock, deep oceans, biogeochemical carbon cycles) continuously for $O(10^4$ to $10^7)$ years. Since the 1970s,

numerical ice sheet models have been developed that are capable of simulating the long-term evolution of kilometers-thick ice cover in three spatial dimensions on regional or continental scales (e.g., Andrews and Mahaffy 1976; Jenssen 1977; Budd and Smith 1979; Oerlemans 1982; Huybrechts 1990; Fastook and Prentice 1994; Greve 1997; Ritz et al. 1997; Verbitsky and Saltzman 1997; Bueller and Brown 2009; Rutt et al. 2009). These models use scaled (approximate) equations of the slow flow of ice deforming under its own weight, and solve for the changing ice thickness and internal ice temperatures on a regular grid, accounting for ice advection, surface snow accumulation minus summer melt, freezing or melting at the ice base, and vertical depression or rebound of the bedrock below the ice. The horizontal grid size in these models is typically 10 to 50 km, with ~ 10 or more vertical layers to keep track of internal temperatures.

Given forcing fields of surface mass balance and sub-ice shelf melting, it is quite feasible to run 3-D ice sheet models over continental domains through 10^4 to 10^7 years, typically taking a few days to weeks of wall-clock time on modern computers. The main impediment to long-term runs is the provision of the forcing fields: net annual mass balance over the ice sheet surface, and rates of oceanic sub-ice-shelf melt, as these vary through the run due to ice-sheet extent and elevations, orbital variations, atmospheric CO_2 , sea level and other forcing. At a minimum two climate fields are needed for the surface budget: seasonal cycles of air temperature and precipitation. Other atmospheric fields such as wind speed and cloudiness can enter into boundary-layer treatments, but their long-term variations tend not to be as important (Pollard and PMIP Participating Groups 2000). The modeling of ocean circulation below ice shelves and sub-ice-shelf melting is challenging and in early stages of development (Beckmann and Goosse 2003; Dinniman et al. 2007; Holland et al. 2008a), and is not discussed here, but could be amenable to some of the temporal methods described below. In the rest of this section, four different approaches are described that provide air temperatures and precipitation over long time scales, roughly in the order they appeared in the literature since the 1970s.

1.1 Energy balance models

The Earth's climate can be crudely approximated as horizontal diffusion of heat, transporting net radiative input from the warm tropics to the cold polar regions, with the transport depending only on the horizontal gradient of a single-level temperature. Ice extent is included implicitly, by increasing solar reflectivity where temperatures are below some value (0 to -5°C). The resulting one-equation Energy Balance Models (EBMs, otherwise known as Simple Climate Models, e.g., Budyko 1969; Sellers 1969; North 1975; cf. Eriksson 1968) are simple enough to be run for millions of years.

In the 1970s, climate models of this kind were coupled with explicit ice-sheet models and applied to the Quaternary ice ages. An even simpler variant is to represent climate by a prescribed snowfall-minus-melt pattern versus height and latitude (Weertman 1976), but the results are very similar if a seasonal EBM is used with ice melt dependent on the EBM's summer temperature (Pollard 1978). In this way, 1-D to 3-D Northern Hemispheric ice-sheet models were driven for 10^5 to 10^6 years though the Quaternary, with the climate component forced primarily by orbital perturbations and sometimes a long gradual cooling representing a presumed

decline in CO₂ (e.g., Oerlemans 1980; Sergin 1980; Birchfield et al. 1982; Pollard 1982, 1983; Hyde and Peltier 1987; Neeman et al. 1988; DeBlonde and Peltier 1993; Mudelsee and Schulz 1997; Tarasov and Peltier 1997, 1999).

More recently, the EBM–ice sheet genre has been applied to earlier glaciations of the Phanerozoic and Precambrian (Hyde et al. 1999, 2000, 2006; Crowley et al. 2001), and to the Cenozoic with climate box-models or inversions (Langebroek et al. 2009; de Boer et al. 2010). An “anomaly” method was introduced, in which only the EBM *differences* from the modern simulated climate are used, and superimposed on modern observed climatology (e.g., Tarasov and Peltier 1997, 1999, for Quaternary cycles). In this method, it is hoped that climate model biases remain the same as for modern, and so cancel to yield more accurate paleoclimate results.

Even using the anomaly method, Energy Balance Models with their simple diffusive heat transport can only crudely capture the largest-scale variations in air temperatures, and offer no information on precipitation changes, or only a little with diffusive moisture-latent heat extensions (Chu and Ledley 1995; Bendtsen 2002; Pollard and Kasting 2005). Although these models have provided important insights and successfully simulated basic ~20 and 40 kyear fluctuations in response to orbital forcing, they leave important questions unanswered, such as the causes of the 100,000 year cycle, Mid-Pleistocene Transition, and rapid initiation of ice sheets at the end of interglacials. Some of the shortcomings might be due to the omission of other long term processes and components (sediment, deep ocean variations, carbon cycling), but they could also be due to oversimplified climate forcing. To test the latter possibility, ice sheet models need to be coupled with more comprehensive climate GCMs or EMICs.

1.2 Basic asynchronous with GCMs and EMICs

For most researchers, the longest full-GCM integrations that are practical for individual experiments are hundreds to several thousand years (Kiehl and Shields 2005; Kahana and Valdes 2009; Liu et al. 2009). For O(10⁶)-year ice-age timescales, continuous GCM integrations through the entire time span are currently infeasible, and GCM applications have been limited to (1) snapshots of individual times, and (2) sub-samples of climate states used creatively to drive the ice sheet model through long periods.

Since the 1980s, many GCM snapshots have been used to deduce mass balance on prescribed North American, Eurasian and Antarctic ice sheets at specific times, mainly ice inception at the end of the last interglacial ~116 ka, the Last Glacial Maximum ~21 ka, and the present to next few hundred years. As well as testing net mass balance, these studies have quantified important forcing and feedback factors such as orbital perturbations, atmospheric *p*CO₂, biosphere–atmosphere and ocean–atmosphere interactions.

The problem of Northern Hemispheric ice-sheet initiation at the end of the last interglacial is amenable to just one climate snapshot, looking at areas with net annual snow accumulation and/or driving nascent ice sheets with an invariant climate, and has received considerable attention (e.g., Royer et al. 1983; Rind et al. 1989; Syktus et al. 1994; Dong and Valdes 1995; DeNoblet et al. 1996; Gallimore and Kutzbach 1996; Pollard and Thompson 1997a; Vettoretti and Peltier 2004).

Many GCM snapshots have also been performed for the Last Glacial Maximum (LGM), looking at the effects of the ice sheets on climate and taking advantage of abundant proxy climate data and reliable ice-sheet reconstructions (e.g., Gates 1976a, b; Manabe and Broccoli 1985; Hall et al. 1996; Fabre et al. 1997; Pollard and Thompson 1997b; Broccoli 2000; Kageyama and Valdes 2000; Hewitt et al. 2003; Kim et al. 2003, 2008; Braconnot et al. 2007). Some of these studies compute the mass balance over the Northern Hemispheric ice sheets, sometimes using the same anomaly method with respect to modern climate as mentioned above for EBMs, and/or interpolating to the finer ice-sheet grids to account for the mismatch in spatial scales as described below (e.g., Charbit et al. 2002, 2007; Zweck and Huybrechts 2005). A few studies have simulated times during the last deglaciation between the LGM and the early Holocene (e.g., Kutzbach and Guetter 1986; Mitchell et al. 1988; Kutzbach et al. 1998; Carlson et al. 2009). Many GCM simulations of modern and future greenhouse climate have been performed, and some GCMs have been used to drive the Greenland and Antarctic ice sheets over the next hundreds to few thousands of years (Huybrechts et al. 2004; Alley et al. 2005; Ridley et al. 2005; Vizcaino et al. 2008).

Deeper-time applications with ice-sheet models driven by one or a few GCM climate solutions include the Cenozoic (DeConto et al. 2007, 2008), Ordovician (Herrmann et al. 2004), Permo-Carboniferous (Horton et al. 2007) and Neoproterozoic Snowball Earth (Baum and Crowley 2001; Donnadieu et al. 2003; Pollard and Kasting 2004).

Since the late 1990s, a few groups have made use of GCMs in longer-term ice-age simulations, using a basic asynchronous method. The ice-sheet model is run continuously, but the GCM is run only a few decades at a time, at several-thousand year intervals using the current orbit and ice sheet size, to provide the mass balance forcing over the ice sheets. The most recently computed mass balance is used to drive the ice sheet model through the next several-thousand-year asynchronous period. Each snapshot requires a few decades of GCM integration to spin up the upper ocean and to average out interannual variability (Pollard et al. 1990). This technique was used by Charbit et al. (2002) to simulate the last Northern Hemispheric deglaciation since 21 ka, and by Horton and Poulsen (2009) in Permo-Carboniferous ice-age experiments. However, the computational expense of the GCM integrations has limited this technique to $O(10^4$ to $10^5)$ year intervals. Another kind of scheme with accelerated orbital variations was used by Jackson and Broccoli (2003) to simulate Northern Hemispheric mass balance in a GCM for the past 165,000 years, but is only applicable for models with no long-term prognostic components, and not for coupling with ice-sheet models.

GCMs with deep-ocean components are needed to test hypotheses concerning interactions with thermohaline circulation and deep-ocean carbon storage (e.g., Imbrie et al. 1992; Toggweiler and Lea 2010). However, this presents a complication for the asynchronous method, because deep-ocean-adjustment time scales are a few thousand years, intermediate between those of the near-surface climate and the ice sheets. Spinning up the deep ocean to equilibrium at each GCM synchronous phase would greatly increase the CPU time required, and may not be realistic. Running the deep ocean using asynchronous and/or acceleration techniques (e.g., Hewitt et al. 2003; Kim et al. 2003) may be feasible, but with more frequent atmospheric updates than for the ice sheets (c.f., Lunt et al. 2006; Timm and Timmermann 2007).

Recently, EMICs (Earth models of intermediate complexity) have been used in basic asynchronous or fully synchronous mode, allowing longer coupled ice-sheet integrations than are possible with GCMs. The atmospheric components in EMICs range from EBMs, statistical-dynamical (zonally averaged eddy) or quasi-geostrophic models to very coarse-resolution GCMs, and are coupled in some cases to relatively sophisticated ocean and other component models (Claussen et al. 2002). Many such studies have been applied to the Quaternary with runs ranging from $\sim 10^4$ to $\sim 10^6$ years, focusing on glacial inception (e.g., Wang and Mysak 2002; Kageyama et al. 2004; Calov et al. 2005, 2009), millennial-scale interactions between ocean thermohaline circulation and ice sheets (e.g., Schmittner et al. 2002; with meltwater experiments intercompared in Rahmstorf et al. 2005), all or part of the last ~ 100 kyear glacial cycle (Gallee et al. 1992; Calov et al. 2002; Charbit et al. 2005; Philippon et al. 2006; Ganopolski et al. 2010), and the last 3 Myear (Berger et al. 1999). They have also been applied to the future (Loutre and Berger 2000; Berger and Loutre 2002; Swingedouw et al. 2008), and deeper-time ice ages (e.g., Donnadiou et al. 2004). However, it is unclear whether the atmospheric dynamics of EMICs are always sufficient for ice-age problems. The representations of atmospheric dynamics and precipitation in some of the simpler EMICs are based empirically on modern climatology, and their skill is likely to decay once they are applied to different boundary conditions, especially over the flanks of past continental ice sheets.

1.3 GCM lookup table

An alternative to the basic asynchronous-GCM method in long-term ice-age experiments is the use of a lookup-table of stored GCM climate solutions. In a preparatory step before long ice-sheet integrations, a collection of GCM snapshots is assembled, forming a look-up table of climates with specified external conditions and ice sizes that span the space of all possible states expected during the long-term runs. Then at any point in a long-term ice sheet run, current temperature and precipitation fields are interpolated appropriately from the appropriate GCM lookup-table members.

The lookup table can consist of just two extreme climate members, often modern and Last Glacial Maximum, that are weighted together in proportion to an empirical “glacial index” such as Greenland ice-core $\delta^{18}\text{O}$ over the last ~ 120 kyears, to drive Northern Hemispheric ice sheets (e.g., Marshall and Clarke 1999a, 2002; Rodgers et al. 2004; Zweck and Huybrechts 2005; Charbit et al. 2007). This is a useful way of investigating processes and testing ice-sheet models, but has the drawback of imposing the Greenland ice-core fluctuations (millennial, orbital and 100,000 yr) directly on the ice-sheet results, so is not ideal for investigating fundamental causes of these variations.

An extension of this method is to build a multi-dimensional matrix of GCM snapshots. Given that (1) there are a limited number of degrees of freedom in the important long-term determinants of climate (orbit, ice-sheet size, and atmospheric CO_2 level), and (2) each determinant is more or less one-dimensional (summer insolation at mid latitudes for orbits, total ice volume for ice-sheet configuration), then a matrix of generic GCM simulations can be assembled with orbits ranging over hot–medium–cold summers, with prescribed ice sheets ranging from none through intermediate to maximum sizes, and with several levels of CO_2 . At any point in a long-term ice-sheet simulation, the climate can be interpolated from members of this

matrix based on the current orbit, overall ice-sheet volume and CO₂ level. A preliminary version of this method was used by DeConto and Pollard (2003) and Pollard and DeConto (2005) to simulate major growth of Cenozoic Antarctic ice at the Eocene–Oligocene boundary ~34 Ma. The full matrix method is currently being applied to Cenozoic Antarctica using a $3 \times 3 \times 3$ matrix with three orbital configurations, three prescribed ice sheet sizes, and three CO₂ levels (Wilson et al. 2009).

However, practical limits on the number of GCM matrix members introduce problems for this method, because orbital configurations and ice-sheet configurations are not really one-dimensional, and need many more than three cases each to fully describe their effects on climate. For instance, the pattern of upslope intensification of precipitation on ice-sheet flanks cannot simply be interpolated from those of two GCM solutions with larger and smaller ice-sheet sizes. Other drawbacks are that (1) other long-term determinants of climate may be important, such as ocean gateway openings, and (2) feedbacks such as meltwater caps and shutdowns of the North Atlantic thermohaline circulation are not readily captured. These drawbacks may be ameliorated by using matrices with more dimensions and elements, but then the number of GCM simulations required may approach or exceed that required by the simpler basic asynchronous method.

1.4 Climate parameterization

Another technique, used from the first long-term ice sheet studies, is simply parameterizing climate over ice sheets. This can be as simple as a zonally symmetric snowfall-minus-snowmelt pattern versus height and latitude, as mentioned in Section 1.1. With two horizontal dimensions, this approach has been limited mostly to Plio-Pleistocene and future Antarctica, with East Antarctic ice sheet (EAIS) configurations similar to present. Modern precipitation and temperature are either taken from climatologic datasets, or parameterized simply in terms of geographic variables. For past times, the modern temperature is perturbed simply in proportion to geologic time series reflecting regional climate, such as Vostok ice core δD or deep-sea-core $\delta^{18}O$ records, plus a lapse-rate correction for elevation changes and a Clausius-Clapeyron relation for precipitation shifts (Huybrechts 1998, 2002; Ritz et al. 2001; Pollard and DeConto 2009). Future projections using this method include Huybrechts and Oerlemans (1990) and Huybrechts et al. (1991), but have been superseded in recent years by GCMs and EMICs (Section 1.2).

The simple parameterizations of modern Antarctic temperature and precipitation are based on regression analyses of modern Antarctic climatology, which relate temperature and precipitation over the EAIS to surface elevation, latitude, surface slope, distance to coast, etc. (Musynski and Birchfield 1985; Fortuin and Oerlemans 1990; Giovinetto et al. 1990). Basic univariate or bivariate regression works quite well for modern EAIS, in part because the ice approximates a single dome centered near the South Pole, and the dependent and independent variables are more or less zonally symmetry. But for earlier Cenozoic ice sheet configurations with non-zonal smaller ice sheets or multiple ice caps, more sophisticated methods would be needed, based possibly on multivariate regression or artificial neural nets (Reusch et al. 2005) and trained by discrete GCM/RCM snapshots of past times. A first step towards this approach in principle was taken by Abe-Ouchi et al. (2007), who forced a Laurentide

ice-sheet model through the last 120 kyears with climate parameterizations adjusted beforehand according to a few discrete GCM snapshots.

2 Multiple spatial scales

2.1 Interpolation from coarse GCM and EMIC grids

The techniques described in the previous section address the *temporal* mismatch between climate and ice sheets. Another problem in all applications using coarse-grid climate models is the mismatch in *spatial* scales. Horizontal grid sizes in GCMs and EMICs are typically several hundred km, and are inadequate to resolve the steep topography around ice-sheet margins that are important in determining ablation. Hence, many studies use straightforward interpolation techniques to derive the mass balance on a much finer-scale ice-sheet model grid.

The first step is to horizontally interpolate the GCM temperature and precipitation fields to the ice-sheet grid, often using simple bilinear interpolation. Simple vertical corrections to air temperature are then applied, to correct for the difference between the GCM-interpolated topography and the actual ice surface elevation, assuming a constant lapse rate (e.g., Thompson and Pollard 1997; Fabre et al. 1998; although it is not clear what this lapse rate should be; see also Greuell et al. 1997). In some cases, a similar vertical correction to precipitation is also applied (elevation-desert effect). If only the monthly means of GCM temperature and other meteorologic variables are saved, then the non-linear effects of diurnal cycles and synoptic variability on melt can be imposed in the calculation of mass balance (see below; Marshall et al. 2004; Zweck and Huybrechts 2005).

Statistical hypsometric corrections to net mass balance for nascent ice caps have been proposed, that account for small-scale topographic variability not even resolved by the ice-sheet grid. This allows for extra melting in deep valleys vs. accumulation on mountain plateaus, but has been used in relatively few long-term studies to date (Walland and Simmonds 1996; Marshall and Clarke 1999b; cf. Kotlarski et al. 2010).

Precipitation is less amenable to simple topographic corrections than air temperature, because it can be affected by forced orographic uplift and descent, resulting for instance in enhanced snowfall on the windward flanks of ice sheets. Air-parcel schemes are available that capture these effects given large-scale incoming winds, temperature, humidity and finer-scale topography (Sanberg and Oerlemans 1983; Fortuin and Oerlemans 1990; Leung and Ghan 1998), but they or simpler versions have been used in only a few ice-sheet studies (Hulton et al. 2002; Calov et al. 2005; van den Berg et al. 2008).

The two main methods used in calculating net annual mass balance at each ice-sheet grid point from seasonal temperature and precipitation are the empirical positive degree days (PDD) parameterization (e.g., Braithwaite 1981; Marshall et al. 2004; Hock 1999 including radiation), and the more physically based surface energy balance model (SEBM; e.g., van de Wal and Oerlemans 1994; Pollard and PMIP Participating Groups 2000; Anslow et al. 2008). Debate continues on the relative merits of these two methods (van de Wal 1996; Bougamont et al. 2007). If not included in the mass-balance model, allowance can be made for partial retention of meltwater by refreezing in the underlying firn, which significantly modifies the

net surface balance (Pfeffer et al. 1991; Thompson and Pollard 1997; Janssens and Huybrechts 2000; Marshall et al. 2004) and can produce a surprising variety of stratified sequences in nature (Shumskii 1964; Paterson 1994). Refreezing is an example of ice-sheet-specific surface physics at the interface between climate and ice sheet models; another example that has not been included in dynamical ice-sheet models to date is the formation and drainage of supraglacial lakes observed on Greenland (Luthje et al. 2006; Box and Ski 2007; Das et al. 2008; Sundai et al. 2009), but such processes are beyond the scope of this paper.

2.2 Regional climate models, and stretched-grid GCMs

Some ice-age problems may require not just smooth interpolation from GCM grids, but explicit modeling of dynamic meteorology over the ice sheets, for instance upslope precipitation and downslope katabatic winds on steep ice-sheet margins. To achieve this, GCM simulations can be enhanced by the use of atmospheric regional climate models (RCMs) on limited domains over a particular ice sheet, driven at their lateral boundaries by GCM meteorology. Several groups have applied RCMs over modern Antarctica and Greenland (e.g., Hines et al. 1997; Bailey and Lynch 2000; Fettweis et al. 2005; van den Broeke et al. 2006; Ettema et al. 2009). Others have used RCMs to investigate details of the mass balance over Northern Hemispheric ice sheets at LGM (Bromwich et al. 2004, 2005) and other times, including the effects of proglacial lakes (Hostetler et al. 2000).

Alternatively, a single GCM can be used with stretched-grid capability, i.e., with a region of finer grid spacing over the area of interest. This technique has been applied to the Eurasian Ice Sheet and proglacial lakes early in the last glacial-interglacial cycle, using the LMDZ GCM (Krinner et al. 2004; Peyaud et al. 2007; see also Colleoni et al. 2009).

In principle, all of the temporal methods described in Section 1 for using GCMs in long-term ice-sheet experiments can be augmented by embedded RCM simulations, by performing one RCM run for each GCM snapshot. This has not yet been attempted, probably due to the computational expense of the RCM (which is comparable to GCMs, given that multiple years of RCM integration would be needed for each snapshot to average out interannual variability).

3 Discussion

In summary, substantial progress has been made since the 1970s in coupled ice sheet–climate modeling techniques, especially in

1. temporal techniques to address the mismatch between time scales (days to years for meteorology and climate, 10^4 to 10^6 years for ice sheets), and
2. spatial techniques to address the mismatch in model grid resolutions (100's km for global climate models, 10's of km for ice-sheet ablation zones).

One aspect of long-term coupled ice sheet–climate simulations that is sometimes overlooked is the need to achieve the correct ice-sheet mass balance not just for one time or interval, but for multiple stages of an ice-age cycle. The interpolation schemes and PDD or SEBM models connecting the GCM climate with ice-sheet mass balance

have uncertain parameters that can be tuned to achieve a desired average mass balance for a particular ice sheet at a particular time (for instance, the Laurentide at LGM). If no such tuning is done, there is wide scatter in whole-ice-sheet mass balances predicted by different GCMs, even though their seasonal climates differ by only a few °C (due to the strong sensitivity of summer melt to air temperatures in ablation zones; Pollard and PMIP Participating Groups 2000). Thus the sign and rate of ice-sheet change can be readily adjusted to a desired value for one particular time, or for a limited unidirectional interval such as the last deglaciation. A much greater challenge is to achieve the right signs and rates of ice-sheet growth and decay through several or all stages of a 100,000 year glacial-interglacial cycle.

The climate and ice-sheet models developed since the 1970s have huge ranges in complexity, and there has been some “vertical” interaction between hierarchical levels (Schneider and Dickinson 1974), mostly for climate models alone. For instance, atmospheric GCM results have been compared and interpreted using single-column radiative–convective models (Henderson-Sellers and McGuffie 1987), EBM ice-age climate sensitivity (Hyde et al. 1989), and conceptual vegetation feedbacks (Brovkin et al. 1998). There have also been some hierarchical studies for ocean models alone (Hirschi and Stocker 2002; Dijkstra and Weijer 2003; cf. Schneider and Thompson 1981).

In the ice sheet–climate arena, hierarchical interactions have mainly been between ice flowline models (with 1 horizontal dimension) and analogous behavior in 3-D models. Lessons learned from 1-D flowline models about Snowball-Earth transitions (Budyko 1969; Sellers 1969; North 1975), Small Ice Cap Instability (SICI, due to albedo feedback; North 1984, Lee and North 1995), Small Ice Sheet Instability and hysteresis (SISI, due to height-mass-balance feedback; Weertman 1961, 1964; Abe-Ouchi and Blatter 1993; Oerlemans 2002), and both SICI and SISI (MacAyeal 1979), helped to interpret tipping points in long-term 3-D model simulations of past and future Greenland and Antarctica (Huybrechts 1993; Crowley and Yip 1994; Crowley and Baum 1995; Cuffey and Marshall 2000; Toniazzo et al. 2004; Pollard and DeConto 2005; Vizcaino et al. 2008) and Neoproterozoic Earth (Donnadieu et al. 2003, 2004; Poulsen 2003; Pollard and Kasting 2004).

Other examples for coupled ice–climate modeling are few. For instance, 0-D conceptual models have explored long-term interactions between global ice volume, temperature and/or carbon dioxide (Saltzman and Verbitsky 1993; Paillard 1995; Parrenin and Paillard 2003), but interactions with 3-D modeling have been sparse, probably due to the gulf in complexity and the difficulty in matching the 0-D model terms with 3-D model processes (although explicit physical processes are identified in Paillard and Parrenin 2004). Similarly, studies using 0-D stochastic resonance models (Benzi et al. 1981; Matteucci 1989) and fundamental analysis of degrees of freedom and free vs. forced variability (Vautard and Ghil 1989; Yiou et al. 1994) have seen little expression in 3-D modeling to date. On the other hand, recent theories and simple models that consider seasonal and interhemispheric aspects of orbital insolation variations and the consequences for ice-sheet cycles (Raymo et al. 2006; Huybers and Denton 2008; Huybers and Tziperman 2008; Huybers 2009) translate more readily to 3-D modeling, and testing them with 3-D ice sheet–climate models will be more straightforward.

Instead of “vertical” interactions in the hierarchical pyramid, much effort in recent years has been directed along the “integration” axis of Claussen et al. (2002),

combining ice sheet–climate models with other long-term components important to ice-sheet evolution, like the evolution of GCMs and EMICs towards comprehensive Earth systems models. These components and processes, with very incomplete modeling references, include:

- Ice stream/shelf flow, grounding-line migration, higher-order flow equations (Weertman 1974; Hughes 1981; MacAyeal 1989; MacAyeal et al. 1996; Hubbard 2000; Pattyn 2002; Bennett 2003 (review); Hindmarsh 2004, 2006, 2009; Payne et al. 2004, 2007; Dupont and Alley 2005; Vieli and Payne 2005; Pollard and DeConto 2007, 2009; Schoof 2007; Bueler and Brown 2009; Hughes 2009; Joughin et al. 2009; Nick et al. 2009)
- Basal hydrology, aquifers, effect on basal stress (Flowers and Clarke 2002; Johnson and Fastook 2002; Parizek and Alley 2004; Clarke 2005 (review); Carlson et al. 2007; Creyts and Schoof (2009)
- Oceanic circulation under ice shelves, oceanic melting/freezing at their base, meltwater/sea-ice feedbacks (Nicholls 1997; Dinniman et al. 2007; Holland et al. 2008a, b; Swingedouw et al. 2008; Thoma et al. 2008; Vizcaino et al. 2008; Walker et al. 2009; Little et al. 2010)
- Ice-shelf and tidewater calving (van der Veen 1996; Doake et al. 1998; Kenneally and Hughes 2002; Benn et al. 2007 (review); Alley et al. 2008; Scambos et al. 2009; Parizek et al. 2010)
- Supraglacial lakes and hydrofracture (Luthje et al. 2006; Box and Ski 2007; Das et al. 2008; Krawczynski et al. 2009; Sundai et al. 2009)
- Subglacial lakes, outburst floods: physical models (Shoemaker 1991; Clarke et al. 2004; Pattyn et al. 2004; Alley et al. 2006; Evatt et al. 2006; Pattyn 2008; Carter et al. 2009)
- Marine bathymetric landforms, interaction with grounding-line advance and retreat (Alley 1991; Anderson 1999; Dahlgren et al. 2002; Alley et al. 2007; Nick et al. 2007)
- Tidal effects on ice streams, stick-slip motion (Bindschadler et al. 2003; Winberry et al. 2009)
- Deformable sediment and till, long-term distribution, effect on basal stress (MacAyeal 1992; Boulton 1996; Alley et al. 1997; Clark and Pollard 1998; Licciardi et al. 1998; Tulaczyk et al. 2000, 2001; Kamb 2001; Bougamont and Tulaczyk 2003; Pollard and DeConto 2007; Hildes et al. 2004)
- Supraglacial debris, stagnation (Vacco et al. 2010)
- Aeolian dust, effects on and of ice sheets (Peltier and Marshall 1995; Overpeck et al. 1996; Mahowald et al. 1999; Krinner et al. 2006; Sugden et al. 2009; Ganopolski et al. 2010)
- Proglacial lakes (Andrews 1973; Pollard 1982; Marshall and Clarke 1999a; Tarasov and Peltier 2005; Peyaud et al. 2007)
- Water isotopic ratios within the ice (Clarke and Marshall 2002; Clarke et al. 2005; Lhomme et al. 2005; Sima et al. 2006)
- Bedrock deformation (Le Meur and Huybrechts 1996; Peltier 2004; Tarasov and Peltier 2004; Lambeck et al. 2006)
- Terrestrial landforms, erosion, landscape evolution (Hindmarsh 1999; Alley et al. 2003; Tomkin 2003; Anderson et al. 2006; Kite and Hindmarsh 2007; Hooke and Fastook 2007; Hindmarsh and Stokes 2008; Jamieson and Sugden

2008; Jamieson et al. 2008; Kessler et al. 2008; Boulton et al. 2009; Sugden 2009 (review); Vacco et al. 2009; Stokes and Tarasov 2010)

These modeling studies provide a rich framework to embed ice-sheet models more realistically into their surrounding environment. The first five to seven on the list are considered to be particularly important for possible fast drawdown through Greenland and West Antarctic ice streams in future decades to centuries, and may be coming into play sooner than was expected a few years ago (e.g., Rignot and Kanagaratnam 2006; Bell 2008; Holland et al. 2008b; Rignot 2008; Vaughan 2008; Pritchard et al. 2009). There is an urgent need to develop more realistic and reliable models of each of them, and much current work in ice sheet–climate modeling is directed to that end.

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